

Periodic flow of heat

When one end of a long metal bar is heated and cooled alternatively then it is called periodic flow of heat . We consider that the temperature varies simple harmonic manner .This method is applied to determine the thermal conductivity of a metal bar .

Fourier equation for one dimensional flow of heat is applied where radiation loss is neglected .

$$\frac{d\theta}{dt} = h \frac{d^2 \theta}{dx^2} \quad (1)$$

Where Diffusivity , $h = \frac{K}{\rho s}$

In order to solve this equation we get

$$\theta = u + v \quad (2)$$

When u and v satisfy the relation

$$d^2 u / dx^2 = 0 \quad (3)$$

and

$$h d^2 v / dx^2 = dv / dt \quad (4)$$

$$h \frac{d^2 v}{dx^2} = \frac{dv}{dt} \quad (4)$$

These equation shows that u is a function and v is a function of x and t both .

The solution of (3) gives the steady temperature at any point on the bar . Thus

$$u = \theta_1 - (\theta_1 - \theta_2 / l) x$$

where θ_1 and θ_2 are mean temperature at the ends of the bar of length l .

Let us now solve (4) ,We have the conditions that

$$v = v_0 e^{\omega t} \quad (5)$$

To solve (4) we can assume a solution of the form

$$\Psi = A e^{i\alpha t + \beta x}$$

Where A , α and β are constants, then

$$\frac{dv}{dt} = i\alpha A e^{i\alpha t + \beta x}$$

$$= i\alpha v$$

$$d^2 v / dx^2 = \beta A e^{i\alpha t + \beta x}$$

$$= \beta^2 v$$

Substituting these values in (4) , we have

$$i \alpha = h \beta^2$$

$$\beta = \pm \sqrt{i\alpha/h}$$

Where $i = \sqrt{-1}$, which can be written as

$$i = \frac{1}{2} (1 + i)^2$$

Then , $\beta = \pm (\alpha / 2h)^{1/2} (1 + i)$

$$A = (\alpha / 2h)^{1/2}$$

and we have

$$v = A e^{\pm a(1+i)x+i\alpha t}$$

The positive sign is physically not allowed, as $v = \infty$ when $x = \infty$.
we take only the negative sign. Thus we get

$$v = A e^{-ax} e^{i(at-ax)} \quad (6)$$

At $x = 0$, this equation reduces to

$$v = A e^{i\alpha t} \quad (7)$$

But we know that at $x = 0$

$$v = v_0 e^{i\alpha t} \quad (8)$$

These two are equivalent , giving

$$A = v_0$$

$$\alpha = \omega$$

Therefore v can be written as

$$v = v_0 e^{-ax} e^{i(\omega t - ax)} \quad (9)$$

considering only the real part , it can be expressed as

$$v = v_0 e^{-ax} \cos (\omega t - ax) \quad (10)$$

In order to make this equation more general , we introduce an arbitrary constant δ in the cosine term . Then we have

$$v = v_0 e^{-ax} \cos (\omega t - ax + \delta) \quad (11)$$

This represents a damped progressive wave . A damped progressive wave having the wave length λ and time period T along x axis is given by

$$v = v_0 e^{-ax} \cos [2\pi (t/T - x/\lambda + \delta)] \quad (12)$$

comparing (11) and (12) , we get

$$\omega = 2\pi / T$$

$$a = 2\pi / \lambda$$

and the velocity of the wave

$$c = \lambda / T = \omega / a$$

where ω is determined from the time period of heating and a is given by

