

Vectors model:

1)  $n \rightarrow 1, 2, 3, \dots$   
 $l \rightarrow 0, 1, 2, \dots$

2) Angular momentum  $l$   
 $l = 0, 1, 2, \dots, (n-1)$   
 $8 \pi a^3 \dots$

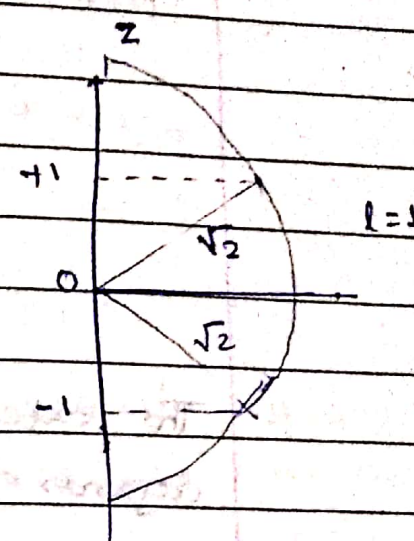
$\Delta l = \pm 1$

3) Space quantization

$\vec{l}$ : vector

magnitude:  $\sqrt{l(l+1)} = l^*$

$\vec{l} = \sqrt{l(l+1)} \hbar$   
 $= l^* \hbar$



It can take values from  $-l$  to  $+l$  and magnetic quantum numbers =  $m_l$ .

Degeneracy =  $(2l+1)$

4) Spin Angular momentum:

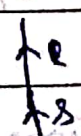
$s = s^* \hbar$   $s = -1/2, +1/2$

$m_s$ : magnetic spin quantum number

Degeneracy for spin =  $2s+1$

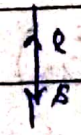
If  $l$  and  $s$  are in one direction

$j = l+s$

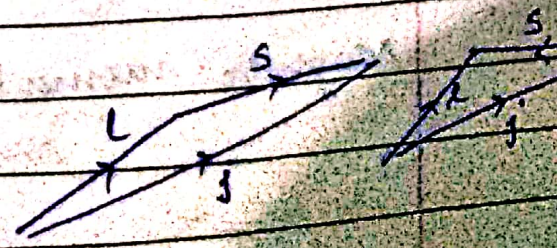


If  $l$  and  $s$  are in opposite "

$j = l-s$



$j$  = total angular momentum



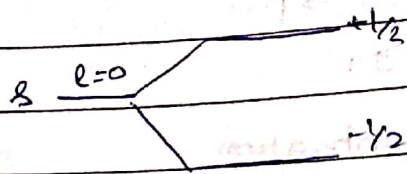
$j = \pm 1/2, \pm 3/2, \dots$

Term for  $l=0$

$2S_{1/2}$

L-S Coupling :  $l = l + s$

The coupling occurs because of orbital angular momentum mag. field arises and the spinning of  $e^-$  also cause a magnetic field and these magnetic fields couple together.

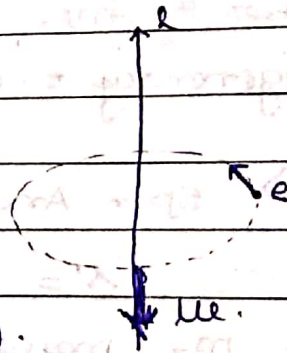


≡ splitting occurs due to coupling and extent of coupling depends on energy separation

The extent of coupling and the separation of energy levels depends on g-factor.

### Bohr Magnetron

The electron in a circular loop produces a current in opposite direction and this leads to magnetic field and magnetic moment.



Current,  $i = e/ET$ . (CGS unit, in MKS unit remove c).

$$T = \text{Time period} = \frac{2\pi r}{v}$$

Then,

$$i = \frac{e v}{2\pi r}$$

Magnetic moment,  $\mu_e = iA$

$$= \frac{e}{c} \left( \frac{v}{2\pi r} \right) \pi r^2$$

$$= \frac{e v r}{2c}$$

$\mu_e$ :  $\perp$  to area but is opp to  $l$ , because  $e^-$  circ

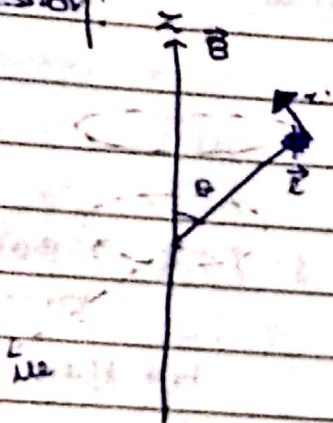
$$\frac{\mu_e}{|r|} = \frac{e \cdot e \cdot \frac{1}{m_e r}}{2c} = \frac{e}{2mc} = \text{independent of orbit}$$

$$\mu_e = \frac{e}{2mc} |r|$$

$$\mu_e = \frac{e}{2mc} (\sqrt{l(l+1)} \hbar)$$

$\frac{\hbar e}{2mc} \equiv$  Independent of the orbit  
 $\equiv$  Bohr magneton  $= 0.928 \times 10^{-20} \text{ erg/cm}$

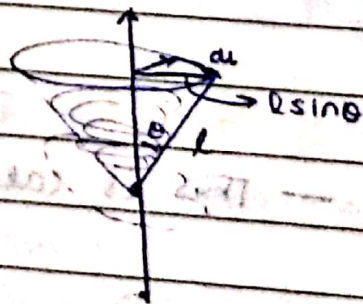
Larmor Precession:



$\vec{l}$  - starts precessing ~~down~~ about magnetic field and the precession has certain frequency called Larmor frequency  $\omega_L$ .

$\theta$ : angle between magnetic field  $B$  and angular momentum vector  $l$ .

$$\mu_e = -\frac{e}{2mc} \vec{l}$$



The interaction with  $B$  produces torque:

$$\vec{\tau} = \vec{\mu}_e \times \vec{B}$$

$\omega \rightarrow$  Angular velocity

$$\omega dt = \frac{|\tau|}{|l| \sin \theta}$$

$$\omega = \frac{|\tau|}{|l| \sin \theta} \frac{1}{\sin \theta} = \frac{\tau}{|l| \sin^2 \theta}$$

$$\tau = \omega |l| \sin \theta$$

$$\omega_L = \frac{|L| \hbar \sin \theta}{I \hbar \sin \theta}$$

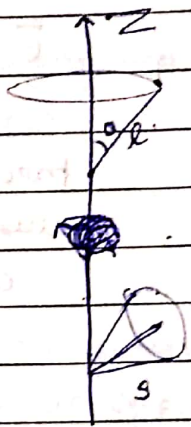
$$= \frac{|L| \hbar}{I}$$

$$\omega_L = \left( \frac{e}{2mc} \right) B$$

B - perturbation due to magnetic field.

if  $B = 10T$ ;

$$\omega_L = \frac{1.606 \times 10^{-19}}{2 \times 9.1 \times 10^{-31} \times 3 \times 10^8} \times 10 = 10^{11} \text{ rad/s}$$



$j = l + s$  : Both of them can precess independently, but there comes a quantity, total angular momentum  $j$  which ~~also~~ remains conserved in always remains in  $z$  direction.

- This is called L-S coupling or spin orbit coupling.

Doublets	Spin orbit coupling & Doublets:
Alkali Metal	$Li \quad 2s^1$
	$Na \quad 3s^1$
	$K \quad 4s^1$
	$Cs \quad 6s^1$

$\equiv$  Single  $e^-$  in outermost orbit

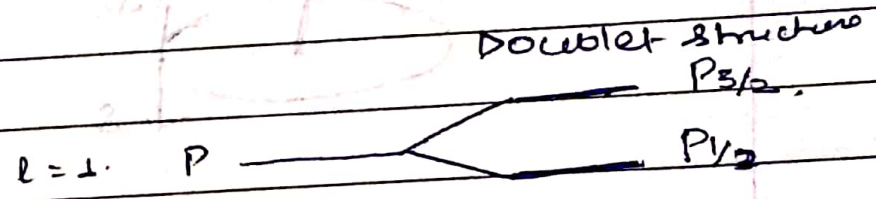
They behave like hydrogen atom because of one  $e^-$  in outermost orbit. Alkali metals always show doublet.

- Doublet separation increases with  $\uparrow$  in atomic number.
- doublet separation in alkali earth metals are larger than corresponding alkali metals.
- Within each element doublet separation  $\downarrow$  on going to higher  $n$ .

Spin orbit interaction.

$l = 0, 1, 2, \dots$   
 $l^* = \sqrt{l(l+1)}$   
 $s^* = \sqrt{s(s+1)}$       $s = +1/2, -1/2$   
 $j^* = \sqrt{j(j+1)}$

Eg:  $l = 1$  ——— P.  
 $j = (l+s) - (l-s)$   
 $= 3/2, 1/2$



*Spin-orbit interaction is more prominent for heavy atoms because of the large number of protons and neutrons in the nucleus.*

$l = 0$

Orbital	$l$	$s$	$j$	Term
s	0	$1/2$	$1/2$	$^2S_{1/2}$
p	1	$1/2$	$1/2, 3/2$	$^2P_{1/2}, ^2P_{3/2}$ → All are doublet structure
d	2	$1/2$	$3/2, 5/2$	$^2D_{3/2}, ^2D_{5/2}$
f	3	$1/2$	$5/2, 7/2$	$F_{5/2}, F_{7/2}$

H  
 Na  
 K  
 Rb  
 Cs

↓  
 Doublet-splitting increases

The fine structure level corresponding to  $j = l - 1/2$  deeper than the corresponding level  $j = l + 1/2$ .

Doublet structure depends on

$z, l, j$

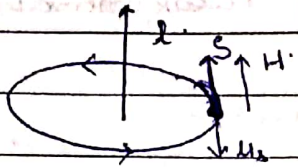
$\propto z$

$\propto 1/l$

Normal order of <sup>doublet</sup> fine structure:

Energy order:  $2P_{3/2} \ 2P_{1/2}$

one configuration



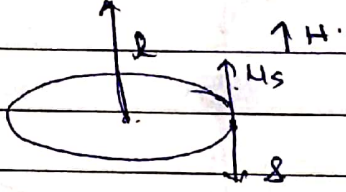
$l = l^* \hbar$

higher energy

$u$  is antiparallel to  $H$  and hence it will tend to align parallel to  $H$ .

$j = l + s$

2nd configuration



$j = l - s$

lower energy

$u \parallel H$

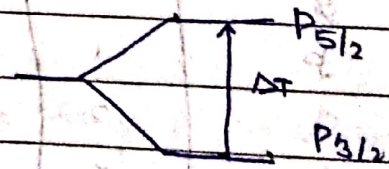
Out of two, the one with lower  $j$ , has lower energy. because classically  $u$  and  $H$  are parallel for lower  $j$  and antiparallel for higher  $j$ .

\* Exception: cesium (Cs)  $7/2$  lies below  $5/2$ .

Splitting - ?

$\Delta T = ?$

$\Delta T \equiv$  depends on strength of interaction

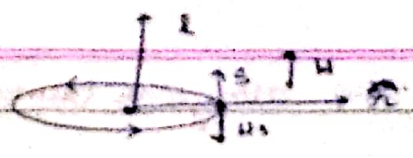


Semi-Classical approach:

$[z, l]$

$l = l^* \hbar \equiv$  angular mom is quantized.

$l^* = \sqrt{l(l+1)}$



Electric field  $\vec{E} = \frac{Ze}{r^3} \vec{r}$

$\vec{H} = \frac{\vec{E} \times \vec{v}}{c}$

$\vec{H} = \frac{Ze}{cr^3} \vec{r} \times \vec{v}$

Bohr Assumption

Bohr quantum:  $2\pi m \vec{r} \times \vec{v} = l^* h$

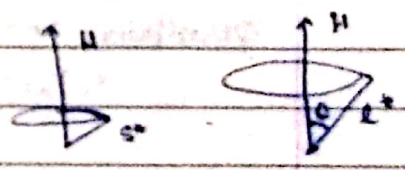
$H = l^* \frac{h}{2\pi} \frac{Ze}{mc} \frac{1}{r^3}$

Under this magnetic field the spin precesses called Larmor precession

Larmor Precession frequency: This for spin motion

$\omega_s = 2\omega_L$

Spin precesses faster than angular momentum around H



Precession  $\equiv$  angle remains constant

Larmor precession frequency for spin motion:  $\omega_L =$  magnetic field  $\times$  ratio of magnetic and mechanical gyromagnetic ratios

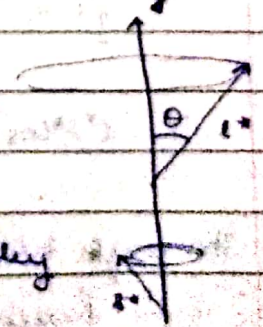
$\omega_L = H \cdot g \cdot \frac{e}{2mc}$

$= l^* \frac{h}{2\pi} \frac{Ze}{mc} \frac{1}{r^3} \frac{2e}{2mc}$

$\omega_L = l^* \frac{h}{2\pi} \frac{1}{r^3} \frac{Ze^2}{mc^2}$

$\rightarrow$  solves the primary Larmor precession

J: conserved



The angle can only be conserved if they have common axis

Interaction Energy

$\Delta U_{int} = \omega_L^* \frac{h}{2\pi} \cos(\theta^*, \theta^*)$

Because of relativistic correction, along with Larmor precession  $\omega_L$  a relativistic precession of half magnitude and opposite direction:

$$\therefore \omega_L' = \omega_L - \omega_r = \text{ordinary precession}$$

$$\therefore \Delta W_{LIS} = \frac{1}{2} \left( \frac{\hbar^2}{2\pi} \frac{1}{a^3} \frac{Ze^2}{m^2 c^2} \right) \cdot \frac{g^* \hbar}{2\pi} \cos(\theta^*, \phi^*)$$

$\equiv$  energy due to coupling of two momenta.

$$\text{or } \Delta W_{LIS} = \frac{\hbar^2}{8\pi^2} \frac{Ze^2}{m^2 c^2} \frac{l^* s^* \cos(\theta^*, \phi^*)}{a^3}$$

$$\langle \Delta W_{LIS} \rangle = \left( \frac{\hbar^2}{8\pi^2} \frac{Ze^2}{m^2 c^2} \right) \langle l^* s^* \cos(\theta^*, \phi^*) \rangle \langle \frac{1}{a^3} \rangle$$

Quantum mechanically,

$$\langle \frac{1}{a^3} \rangle = \frac{Z^3}{a_0^3 n^3}$$

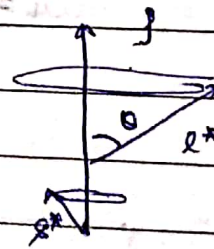
$$\frac{1}{a_0^3 n^3} \frac{1}{(l+1)(l+1/2)}$$

$a_0$ : Radius of 1st orbit

$l \rightarrow$  angular quantum no.

$n \rightarrow$  principal q.n.

angle b/w  $l^*$  &  $s^*$



from vector diagram

$$j^2 = l^2 + s^2 + 2l^* s^* \cos(\theta^*, \phi^*)$$

$$l^* s^* \cos(\theta^*, \phi^*) = \frac{j(j+1) - s(s+1) - l(l+1)}{2}$$

The angle b/w  $l^*$  and  $s^*$  remains conserved all the time and hence average is same as the usual value.

$$\Delta W_{LIS} = \frac{Ze^2}{2m^2 c^2} \frac{\hbar^3}{4\pi^2} \frac{Z^3}{a_0^3 n^3 (l+1)(l+1/2)} \left( \frac{j(j+1) - s(s+1) - l(l+1)}{2} \right)$$



Rydberg constant  $R = \frac{2\pi me^4}{ch^3}$

fine structure constant,  $\alpha = \frac{4\pi^2 e^4}{c^2 h^3}$

$$\Delta \omega_{e,s} = \left( \frac{R \alpha^2 z^4 ch}{n^3 l(l+1)(l+1/2)} \right) \left( \frac{j(j+1) - l(l+1) - s(s+1)}{2} \right)$$

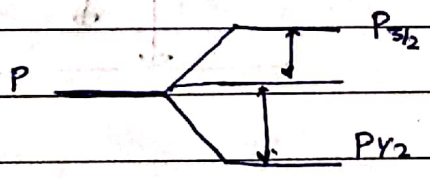
or  $\frac{\Delta \omega_{e,s}}{hc} = a \frac{j(j+1) - l(l+1) - s(s+1)}{2}$

where  $a = \frac{R \alpha^2 z^4}{n^3 l(l+1)(l+1/2)}$

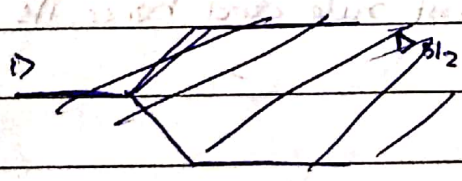
$$\Delta T = \frac{\Delta \bar{\omega}_{e,s}}{hc} = a \frac{j(j+1) - l(l+1) - s(s+1)}{2} = -\Gamma$$

Term T =  $T_0 - \Gamma$   
 ↓  
 Centre of term

$\Delta T < z \therefore z \uparrow \equiv \text{separation} \uparrow$



- The splittings are asymmetric because of different j values.
- Splitting ↓ as l ↑.



For P term:

$$\Delta \bar{\omega} \Big|_{P_{3/2}} = a \frac{3/2 \times 5/2 - 2 - 1/2 \times 1/2}{2} = a/2$$

$$\Delta \bar{\omega} \Big|_{P_{1/2}} = 3a/2$$