

Let the coordinates of the radius of the orbit are  $x, y, z$  . Then we have

$$a^2 = x^2 + y^2 + z^2$$

If  $Z$  – axis is chosen to be in the plane perpendicular to the applied field , then

$$a_1^2 = x^2 + y^2$$

The mean value of all possible radius is given by

$$a^2 = x^2 + y^2 + z^2 = 3x^2$$

$$= \frac{2}{3} a^2 \quad (13)$$

Using (13) in (12) ,

$$\Delta M = -\left( \frac{e^2 H}{4m} \right) \left( \frac{2}{3} a^2 \right)$$

$$\Delta M = -\left( \frac{e^2 H}{m} \right) \left( \frac{1}{6} a^2 \right)$$

For an atomic number  $Z$  , there are  $Z$  – electronic orbits and there orbits may lie in all possible directions in space , hence total induced magnetic moment .

$$M = -e^2 H / 6m \sum a^2 \quad (14)$$

Let  $N$  = number of atoms per unit volume

Then the magnetic moment per unit volume or intensity of magnetisation ,

$$I = -e^2 H N / 6m \sum a^2$$

Hence the diamagnetic susceptibility

$$\chi = I / H = -e^2 N / 6m \sum a^2 \quad (15)$$

This relation is independent of the temperature and field strength which agrees with experimental result .

This formulae is endorsed by quantum mechanics and the temperature - independency of this value may be understood from the evaluation of  $\sum a^2$  using the wave functions.

The energy difference between the wave functions with different radial parts is approximates 10 eV.

The thermal energy  $kT$  at room temperature amounts to approximately 1/40 eV. Thus, diamagnetism exhibits no dependence on temperature.

Atomic diamagnetism usually is of the order of  $\sim 10^{-6}$  emu, and increases in absolute magnitude for larger atoms with a bigger atomic number because they have a wider radial distribution function.