

Application of Hamiltonian to the problem of motion of a charged particle in an electromagnetic field.

The force on a particle with charge q is given by Lorentz formula

$$F = q \left[\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right] \quad \text{--- (1)}$$

where the electric and magnetic field in vacuum can be expressed in the form

$$\vec{B} = \nabla \times \vec{A} \quad \text{and} \quad E = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t},$$

where \vec{A} is the vector potential and ϕ the scalar potential, c is the velocity of light in vacuum.

The equation of motion is given by

$$\begin{aligned} \vec{F} = \frac{d}{dt}(m\vec{v}) &= q \left[\vec{E} + \frac{\vec{v}}{c} \times \vec{B} \right] \\ &= q \left[-\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} + \frac{\vec{v}}{c} \times (\nabla \times \vec{A}) \right] \quad \text{--- (2)} \end{aligned}$$

We know that $\vec{v} \times (\nabla \times \vec{A}) = \nabla(\vec{v} \cdot \vec{A}) - (\nabla \cdot \vec{v}) \vec{A}$

thus we have

$$\begin{aligned} \frac{d}{dt}(m\vec{v}) &= q \left[-\nabla \phi - \frac{1}{c} \left(\frac{\partial \vec{A}}{\partial t} + \vec{v} \cdot \nabla \vec{A} \right) + \nabla \frac{\vec{v} \cdot \vec{A}}{c} \right] \\ &= e \left[-\nabla \left(\phi - \frac{\vec{v} \cdot \vec{A}}{c} \right) - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right] \end{aligned}$$

$$\text{or} \quad \frac{d}{dt} \left(m\vec{v} + \frac{q}{c} \vec{A} \right) = -q \nabla \left(\phi - \frac{\vec{v} \cdot \vec{A}}{c} \right) \quad \text{--- (3)}$$

This equation has the general form of a set of Lagrange's equations

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) = \frac{\partial L}{\partial q_k}$$

$$\text{Here} \quad \left. \begin{aligned} \frac{\partial L}{\partial \dot{q}_k} &= \frac{\partial L}{\partial \dot{x}_i} = m \dot{x}_i + \frac{q}{c} A_i \\ \frac{\partial L}{\partial q_k} &= \frac{\partial L}{\partial x_i} = \frac{\partial}{\partial x_i} \left[-q \left(\phi - \frac{v_j A_j}{c} \right) \right] \end{aligned} \right\} \quad \text{--- (4)}$$

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or in vector notation

$$\frac{\partial L}{\partial \vec{v}} = m\vec{v} + \frac{q}{c}\vec{A} \quad \text{--- (5)}$$

$$\& \quad \frac{\partial L}{\partial x} = \frac{\partial}{\partial x} \left\{ -q\left(\phi - \frac{\vec{v} \cdot \vec{A}}{c}\right) \right\}$$

These relations (4) & (5) are true only when

$$L = \frac{1}{2}mv^2 - q\phi + q \frac{\vec{v} \cdot \vec{A}}{c}$$

$$\text{Momentum } \vec{P} = \frac{\partial L}{\partial \vec{v}} = m\vec{v} + \frac{q}{c}\vec{A} \quad \text{--- (6)}$$

Thus the momentum of a charged particle moving in electromagnetic field is not defined as its mass times velocity.

From eqn. (6), we may write

$$\vec{v} = \frac{1}{m} \left[\vec{P} - \frac{q}{c}\vec{A} \right]$$

Hence the Hamiltonian H is given by

$$H = \frac{1}{2m} \left[\vec{P} - \frac{q}{c}\vec{A}(\vec{r}, t) \right]^2 + q\phi(\vec{r}, t) \quad \text{--- (7)}$$

The canonical equations are then (7), together with

$$\vec{P} = -\vec{\nabla}H = -q\vec{\nabla}\phi + \frac{q}{c}\vec{\nabla}(\vec{v} \cdot \vec{A}) \quad \text{--- (8)}$$

This is same as Lorentz force equation if we write kinetic energy T

$$T = \frac{1}{2}mv^2 = E - q\phi$$

we can write

$$\vec{P}_k = m\vec{v}_k = \vec{P}_k - \frac{q}{c}\vec{A}_k$$

For the kinetic momentum as the difference between the total momentum and Em. momentum the kinetic momentum is, of course, related to $\vec{k} \cdot \vec{E}$ &

$$T = \frac{1}{2m} \sum_k P_k^2 = \frac{P^2}{2m} \text{ but is not canonically conjugate.}$$

i.e. $P_k \neq \frac{\partial H}{\partial x_k}, x_k \neq \frac{\partial H}{\partial P_k}$

The Lagrangian (nonrelativistic) for a single particle

$$L = \frac{1}{2} m v^2 - q\phi + \frac{q}{c} \vec{A} \cdot \vec{v}$$

with canonical momenta

$$p_i = m v_i + \frac{q}{c} A_i$$

From the definition it is given by

$$H = \sum_i p_i v_i - L \\ = m v^2 + \frac{q}{c} \vec{A} \cdot \vec{v} - L$$

or finally $H = \frac{m v^2}{2} + q\phi = T + q\phi$ the total energy of the particle.

In terms of momenta, the Hamiltonian appears as

$$H = \frac{1}{2m} \left(p - \frac{q}{c} A \right)^2 + q\phi$$

A similar result is obtained for the relativistic Hamiltonian with electromagnetic forces. Here too the canonical momenta ($p_i = m v_i + \frac{q}{c} A_i$) have the additional expression $\frac{q A_i}{c}$ which is just such as to eliminate the term in Hamiltonian involving vector potential. The H is again

$$H = T + q\phi$$

as in the relativistic case

$$H = \sqrt{\left(\vec{p} - \frac{q \vec{A}}{c} \right)^2 c^2 + m^2 c^4} + q\phi$$

The End