

Variational or Hamilton's principle.

A more general formulation of the Lagrangian in mechanics is due to the variational principle (so called Hamilton's principle). The principle is stated in a generalised form independent of any co-ordinates system and hence is useful in non-mechanical systems and fields also. It is also known as integral principle.

This principle may be stated as — "out of all the possible paths along which a dynamical system may move from one point to another point within a given interval of time (consistent with constraints, if any), the actual path followed is that which minimizes the time integral of the Lagrangian".

Analytically it can be represented as

$$I = \int_{t_1}^{t_2} L dt = \text{extremum} \quad \dots \quad (1)$$

where I is the extremum value of time integral of Lagrangian and is known as Hamilton's principal function of the path.

Taking δ -variation of the eqn. (1), the variational principle may also be represented as

$$\delta I = \delta \int_{t_1}^{t_2} L dt = 0 \quad \dots \quad (2)$$

where $L = T - V$, $T = K.E$ & $V = P.E.$

Ques. State the principle of least action and derive the equations of motion.

Deduction of Hamilton's Principle

Let us consider that the conservative holonomic dynamical system moves from P to Q, where P & Q are initial and final points at time t_1 & t_2 respectively. Let PRQ be the actual path and PR'Q & PR''Q be the two neighbouring paths out of infinite no. of possibilities.

For the deductions of this principle there must be satisfied two following conditions -

δt must be zero at end points
and $\delta \dot{x}$ must be zero at end points.

Let the ^{ith particle of the} system be acted upon by a no. of forces given by \vec{F}_i acquiring acceleration $\ddot{\vec{x}}_i$, so that we have

$$\vec{F}_i = m_i \ddot{\vec{x}}_i$$

From D'Alembert's Principle, we have

$$\sum (\vec{F}_i - m_i \ddot{\vec{x}}_i) \delta \dot{x}_i = 0$$

$$\text{or } \sum_i \vec{F}_i \cdot \delta \dot{x}_i - \sum_i m_i \ddot{\vec{x}}_i \cdot \delta \dot{x}_i = 0 \quad \text{--- (1)}$$

$$\text{But } \ddot{\vec{x}}_i \cdot \delta \dot{x}_i = \frac{d}{dt} (\dot{\vec{x}}_i \cdot \delta \dot{x}_i) - \dot{\vec{x}}_i \cdot \frac{d}{dt} (\delta \dot{x}_i) \quad \text{--- (2)}$$

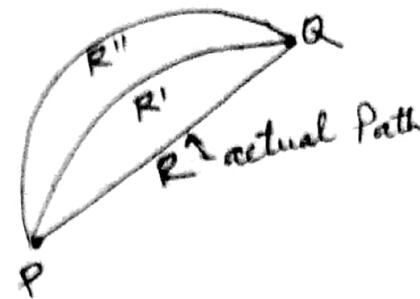
If there is a little variation along the actual and neighbouring paths, we have

$$\frac{d}{dt} (\delta \dot{x}_i) = \delta \frac{d}{dt} (\dot{\vec{x}}_i) = \delta \dot{\vec{x}}_i \quad \text{--- (3)}$$

using eqn (3), eqn (2) may be written as

$$\ddot{\vec{x}}_i \cdot \delta \dot{x}_i = \frac{d}{dt} (\dot{\vec{x}}_i \cdot \delta \dot{x}_i) - \dot{\vec{x}}_i \cdot \delta \dot{\vec{x}}_i \quad \text{--- (4)}$$

using eqn (4), eqn (1) becomes



$$\sum_i \underline{F}_i \cdot \delta \underline{r}_i - \sum_i m_i \left[\frac{d}{dt} (\dot{\underline{r}}_i \cdot \delta \underline{r}_i) - \dot{\underline{r}}_i \cdot \delta \dot{\underline{r}}_i \right] = 0$$

or $\sum_i \underline{F}_i \cdot \delta \underline{r}_i - \sum_i m_i \left[\frac{d}{dt} (\dot{\underline{r}}_i \cdot \delta \underline{r}_i) - \frac{1}{2} \delta (\dot{\underline{r}}_i^2) \right] = 0$

or $\sum_i \underline{F}_i \cdot \delta \underline{r}_i + \sum_i \frac{1}{2} m_i \delta (\dot{\underline{r}}_i^2) = \sum_i \frac{d}{dt} (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i)$

or $\sum_i \underline{F}_i \cdot \delta \underline{r}_i + \delta \sum_i \left(\frac{1}{2} m_i \dot{\underline{r}}_i^2 \right) = \sum_i \frac{d}{dt} (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i) \rightarrow (5)$

But $\sum_i \underline{F}_i \cdot \delta \underline{r}_i = \delta W = \text{work done by the forces } \underline{F}_i \text{ during displacement}$
 $\delta \underline{r}_i$

and $\delta \sum_i \left(\frac{1}{2} m_i \dot{\underline{r}}_i^2 \right) = \delta T$ where T is K.E

Therefore eqn. (5) becomes

$$\delta W + \delta T = \sum_i \frac{d}{dt} (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i)$$

Integrating above eqn between limits t_1 & t_2 , we get

$$\begin{aligned} \int_{t_1}^{t_2} (\delta W + \delta T) dt &= \int_{t_1}^{t_2} \sum_i \frac{d}{dt} (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i) dt \\ &= \sum_i \int_{t_1}^{t_2} d (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i) \\ &= \sum_i (m_i \dot{\underline{r}}_i \cdot \delta \underline{r}_i) \Big|_P^Q = 0 \end{aligned}$$

Since $\delta \underline{r}_i$ at P & Q is zero

For a conservative system, we know that

$\delta W = -\delta V$ where V is potential energy

$$\therefore \int_{t_1}^{t_2} (-\delta V + \delta T) dt = 0 \quad \text{or} \quad \int_{t_1}^{t_2} \delta(T - V) dt = 0$$

$$\text{or } \delta \int_{t_1}^{t_2} L dt = 0 \quad \text{or} \quad \int_{t_1}^{t_2} L dt = \text{extremum}$$

which is Hamilton's principle